

## Lecture 23

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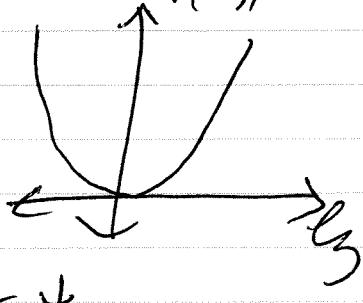
### Hermite & Laguerre Polynomials

Consider a 1D quantum simple harmonic oscillator:

$$-\frac{\hbar^2}{2m} \psi''(\xi) + V(\xi) \psi(\xi) = E \psi(\xi)$$

$$V(\xi) = \frac{1}{2} m \omega^2 \xi^2 \quad (\text{quadratic potential})$$

$\omega$  - frequency



$$\Rightarrow -\frac{\hbar^2}{2m} \psi'' + \frac{1}{2} m \omega^2 \xi^2 \psi = E \psi$$

$$\Rightarrow \psi'' - \frac{m^2 \omega^2}{\hbar^2} \xi^2 \psi + \frac{2m}{\hbar^2} \psi = 0$$

$\xi$  has units of length  
 $m \omega \hbar^{-1}$  has units ~~length~~ [length]<sup>2</sup>

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Then  $x = \xi \sqrt{\frac{mw}{\hbar}}$  is dimensionless

Define  $\ell = \sqrt{\frac{\hbar}{mw}}$

& rescale  $x = \xi/\ell$  &  $\frac{dx}{d\xi} = \frac{1}{\ell}$   
 $\psi(\xi) = y(x)$

$$\Rightarrow \frac{d\psi}{d\xi} = \frac{dx}{d\xi} \frac{dy}{dx} = \frac{1}{\ell} \frac{dy}{dx}$$

$$\Rightarrow \psi''(\xi) = \frac{1}{\ell^2} y''(x)$$

$$\Rightarrow \frac{1}{\ell^2} y''(x) - \frac{1}{\ell^2} x^2 y(x) = -\frac{2mE}{\hbar^2} y(x)$$

$$\Rightarrow y''(x) - x^2 y(x) = -\frac{2mE}{\hbar^2} \frac{\hbar}{mw}$$

$$= -\frac{2E}{\hbar^2 w}$$

define  $\varepsilon = \frac{E}{\hbar^2 w}$

(3)

$$\Rightarrow y''(x) - x^2 y(x) = -2\epsilon y(x)$$

- we do not justify, but ~~do~~ state that energy is quantized, so that  $\epsilon = n + 1/2$  for  $n \in \{0, 1, 2, \dots\}$
- we then have

$$y_n''(x) - x^2 y_n(x) = -2(n + 1/2)y_n(x)$$

for  $n \in \{0, 1, 2, \dots\}$

This is a Hermite diff eq.

where  $y_n(x)$  are Hermite functions.

- Let  $\frac{d}{dx} = D = D_x = \partial_x = \delta$

$$\begin{aligned} \text{then } [D-x][D+x]y &= [D-x][y'(x) + xy] \\ &= y''(x) + y + xy'(x) - xy'(x) - x^2 y \end{aligned}$$

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$$= y''(x) + y - x^2y$$

Furthermore,

$$[D+x][D-x]y = [D+x](y' - xy)$$

$$\begin{aligned} &= y'' - y - xy' + xy' - x^2y \\ &= y'' - y - x^2y \end{aligned}$$

Using this, we can write

Hermite diff. eq. as

$$y_n''(x) - x^2y_n(x) + y_n(x) = -2ny_n(x)$$

$$(1) \Rightarrow [D-x][D+x]y_n(x) = -2ny_n(x)$$

can also write as

$$y_n''(x) - x^2y_n(x) - y_n(x) = -2(n+1)y_n(x)$$

$$(2) \Rightarrow [D+x][D-x]y_n(x) = -2(n+1)y_n(x)$$

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Now operate on 1st one w/  $D+x$

& on second one w/  $D-x$

+ change n to m

$$(3) \Rightarrow [D+x][D-x][D+x] y_m(x) = [D+x]^{2m} y_m(x)$$

$$(4) \Rightarrow [D-x][D+x][D-x] y_m(x) = -2(m+1) [D-x] y_m(x)$$

Compare (1) & (4).

If  $y_n(x) = (D-x)y_m(x)$  &

$n=m+1$ , then

the equations are identical.

Then write

$$y_{m+1}(x) = (D-x)(y_m(x))$$

& can conclude from this development,  
that if we have a solution  $y_m(x)$   
for  $n=m$ , we can obtain

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another solution for  $n=m+1$

by applying the "raising operator"

$$[D-x] \rightarrow y_{m+1}(x).$$

- Similarly, from (2) + (3),  
we find that

$$y_{m-1}(x) = [D+x] y_m(x)$$

$[D+x]$  is known as a "lowering operator"

- these operators are also called ladder operators since they all allow for going up or down, as in a rung of ladders.

- Idea is now to solve Hermite diff. eq. for  $n=0$  + then generate all other solutions by using raising operator.

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$$y_{m+1}(x) = [D - x] y_m(x)$$

Consider  $n=0$ . Then equation is

$$y_0'' - x^2 y_0 = -y_0$$

$$\Rightarrow y_0'' + (1-x^2)y_0 = 0$$

$$\text{Consider } y_0(x) = e^{-x^2/2}$$

$$\begin{aligned} \text{take } y_0'(x) &= e^{-x^2/2} \cdot (-x) \\ &= -x y_0(x) \end{aligned}$$

$$\begin{aligned} y_0''(x) &= -y_0(x) + x^2 y_0(x) \\ &= -(1-x^2) y_0(x) \end{aligned}$$

$\Rightarrow y_0(x)$  is a solution to

$$y_0''(x) + (1-x^2) y_0(x) = 0.$$

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So then the solutions are given by

$$y_n(x) = [D - x J^n] e^{-x^2/2}$$

(called Hermite functions) for  $n \in \{0, 1, 2, \dots\}$ .

can also write this <sup>more simply</sup> as

$$y_n(x) = e^{x^2/2} \left[ \frac{d}{dx} \right]^n e^{-x^2}$$

Proof: Consider that

$$e^{x^2/2} D \left[ e^{-x^2/2} f \right]$$

$$= e^{x^2/2} \left[ -x e^{-x^2/2} f + e^{-x^2/2} f' \right]$$

$$= -x f + f'$$

$$= [D - x] f$$

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So then for the base step, we get

$$\begin{aligned}
 & [D-x]^2 e^{-x^2/2} = \cancel{e^{x^2/2}} \\
 & [D-x] [D-x] (e^{-x^2/2}) \\
 & = [D-x] \left( e^{x^2/2} D [e^{-x^2}] \right) \\
 & = e^{x^2/2} D \left[ e^{-x^2/2} \left( e^{x^2/2} D [e^{-x^2}] \right) \right] \\
 & = e^{x^2/2} D D e^{-x^2} \\
 & = e^{x^2/2} D^2 e^{-x^2}
 \end{aligned}$$

Now for inductive step:

Suppose that

$$[D-x]^n e^{-x^2/2} = e^{x^2/2} D^n e^{-x^2}$$

Then

$$[D-x]^{n+1} e^{-x^2/2} = [D-x] [D-x]^n e^{-x^2}$$

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$$= [D - x] \left( e^{x^2/2} D^n e^{-x^2} \right)$$

$$= e^{x^2/2} D \left[ e^{-x^2/2} \left( e^{x^2/2} D^n e^{-x^2} \right) \right]$$

$$= e^{x^2/2} D \circ D^n e^{-x^2}$$

$$= e^{x^2/2} D^{n+1} e^{-x^2}$$

Q.E.D.

multiplying Hermite functions by

$(-1)^n e^{x^2/2}$  gives Hermite polynomials, defined as

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}$$

formula always gives rise to  
a polynomial of degree  $n$ .

$$H_0(x) = (-1)^0 e^{x^2} D^0 e^{-x^2} = 1$$

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$$\begin{aligned}
 H_1(x) &= (-1)^1 e^{x^2} D^1 e^{-x^2} \\
 &= (-1) e^{x^2} (-2x) e^{-x^2} \\
 &= 2x
 \end{aligned}$$

$$\begin{aligned}
 H_2(x) &= (-1)^2 e^{x^2} D^2 e^{-x^2} \\
 &= (-1)^2 e^{x^2} D \left\{ -2x e^{-x^2} \right\} \\
 &= e^{x^2} \left[ -2e^{-x^2} + 4x^2 e^{-x^2} \right] \\
 &= 4x^2 - 2
 \end{aligned}$$

Hermite polynomials satisfy  
the Hermite equation

$$y'' - 2xy' + 2n y = 0$$

They are also orthogonal on  
 $(-\infty, \infty)$  w.r.t weight function  $e^{-x^2}$

$$\int_{-\infty}^{\infty} dx e^{-x^2} H_n(x) H_m(x) = \delta_{nm} \sqrt{\pi} 2^n n!$$

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The solution to the original  
Hermite diff. eq. 13

$$y_n(x) = (-1)^n e^{-x^2/2} H_n(x)$$

Putting in units of length, we get

$$\psi_n(\ell \xi) = y_n(x) \cdot c_n$$

$$\Rightarrow \psi_n(\ell \xi) = c_n (-1)^n e^{-(\xi/\ell)^2/2} H_n(\xi/\ell)$$

↑ normalization

where  $\ell = \sqrt{\frac{\hbar}{m w}}$  &  $c_n$  is a norm.

To compute  $c_n$ , calculate

$$1 = \int_{-\infty}^{\infty} |\psi_n(\ell \xi)|^2 d\xi$$

$$= |c_n|^2 \ell \int_{-\infty}^{\infty} e^{-(\xi/\ell)^2/2} H_n^2(\xi/\ell) \frac{d\xi}{dQ}$$

(B)

$$= |c_n|^2 \ell \int_{-\infty}^{\infty} e^{-x^2} H_n^2(x) dx$$

$$= |c_n|^2 \ell \sqrt{\pi} 2^n n!$$

$$\Rightarrow c_n = [\ell \sqrt{\pi} 2^n n!]^{-1/2}$$
$$= \left(\frac{mw}{\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}}$$

So then

$$\psi_n(\ell) = \left(\frac{mw}{\hbar}\right)^{1/4} \frac{(-1)^n}{\sqrt{2^n n!}} e^{-\frac{mw}{2\hbar} \frac{x^2}{3}} H_n\left(3\sqrt{\frac{mw}{\hbar}}\right)$$

choose units s.t.  $\ell=1$  & plot

$$\Psi_n = \epsilon_n + \psi_n$$

