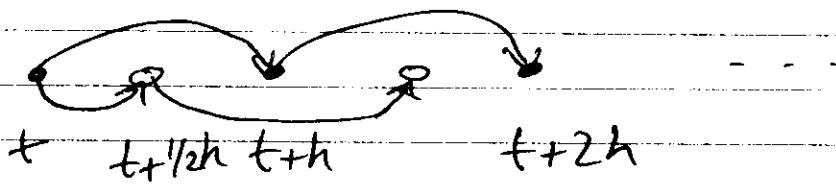


## Lecture 18

Recall leap-frog method



Equations were

$$(*) \quad \begin{aligned} x(t+h) &= x(t) + h f(x(t+1/2h), t+1/2h) \\ x(t+3/2h) &= x(t+1/2h) + h f(x(t+h), t+h) \end{aligned}$$

The method is time reversal  
symmetric

Substitute  $h \rightarrow -h$  & get

$$x(t-h) = x(t) - h f(x(t-1/2h), t-1/2h)$$

$$x(t-3/2h) = x(t-1/2h) - h f(x(t-h), t-h)$$

Now time shift  $t \rightarrow t+3/2h$   
& get

(2)

$$x(t+1/2h) = x(t+3/2h) - h f(x(t+h), t+h)$$

$$x(t) = x(t+h) - h f(x(t+1/2h), t+1/2h)$$

Compare to (\*) & we see that  
 the equations are the same,  
 but run backwards

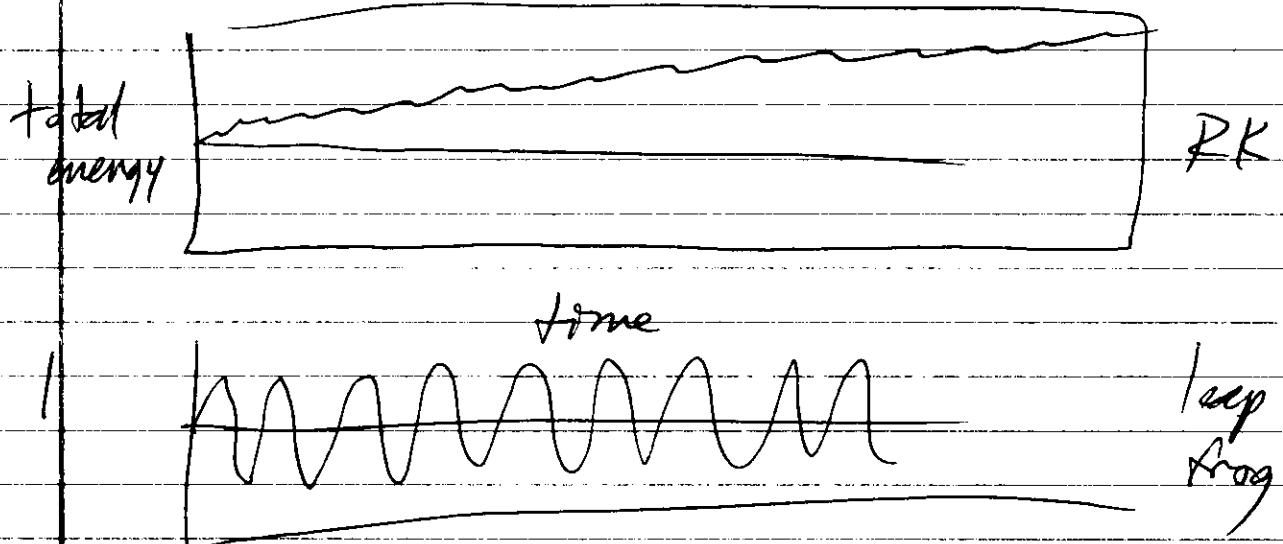
This is not true of Runge-Kutta  
 (can do the same calculation &  
 you won't have the time  
 reversal symmetry)

Why is the time reversal symmetry  
 important? One reason is  
 conservation of energy.

When using Runge-Kutta for  
 nonlinear pendulum, the total  
 energy fluctuates & drifts over time

(3)

For the leap-frog method,  
the total energy fluctuates but  
there is no drift



So leap-frog method is useful

for solving energy conserving  
physical systems over long periods  
of time.

(4)

## Verlet method

variation of leap frog

Suppose equations of motion take  
the form

$$\frac{d^2x}{dt^2} = f(x, t)$$

e.g.,  $F = ma$

can convert to

$$\frac{dx}{dt} = v \quad \frac{dv}{dt} = f(x, t)$$

could directly apply leap frog  
method to vector  $\underline{r} = (x, v)$

w/  $\frac{d\underline{r}}{dt} = \underline{f}(\underline{r}, t)$

But let's instead write out  
the leap frog method in full

$$x(t+h) = x(t) + h v(t+\frac{1}{2}h) \quad (1^{st} q.)$$

$$v(t+\frac{3}{2}h) = v(t+\frac{1}{2}h) + h f(x(t+h), t+h)$$

We can derive a full solution using these two equations alone. The vector method requires double the work, but here we only calculate  $x @ t + kh$   
 $+ v @ t + (1/2 + k)h$

- works for diff eq's w/ special

form

$$\frac{dx}{dt} = v$$

↑  
doesn't  
depend on  
 $x$

$$\frac{dv}{dt} = f(x, t)$$

↑  
doesn't  
depend  
on  $v$

- Potential issue: What if we want to calculate quantity that depends on  $x$  or  $v$ , like total energy?

we only have  $x @ t + kh$

$$+ v @ t + (k + 1/2)h$$

(6)

Supposing that we did know  
 $v(t+h)$  then we could calculate  ~~$v(t+1/2h)$~~   
 ~~$v(t+1/2h)$~~   $v(t+1/2h)$  by going  
 backwards w/ the Euler method  
 giving

$$v(t+1/2h) = v(t+h) - \frac{1}{2}h f(x(t+h), t+h)$$

then rearrange as

$$v(t+h) = v(t+1/2h) + \frac{1}{2}h f(x(t+h), t+h)$$

So the full Verlet method is

$$x(t+h) = x(t) + h v(t+1/2h)$$

$$k = h f(x(t+h), t+h)$$

$$v(t+h) = v(t+1/2h) + \frac{1}{2}k$$

$$v(t+3/2h) = v(t+1/2h) + k$$

put "bars" underneath to get  
 vector quantities...

(7)

## Modified midpoint method

Another advantage of leap frog method is that the total error is an even function of the step size  $h$  (due to time reversal symmetry)

$\Rightarrow$  expansion of error in power series of  $h$  has only even terms & no odd terms.

We can see this in more detail :

A single step of leap frog method

is accurate to order  $h^2$  & has error to order  $h^3$

Write error as  $\epsilon(h)$  w/

first term proportional to  $h^3$ .

What do the other terms look like?

(8)

Take a small step<sup>^</sup> w/ leap  
frog method.

Gives the solution plus  $\varepsilon(h)$   
(error)

Now go backwards, i.e., step size is  $-h$

Due to time symmetry,

the change in the solution is  
the reverse of the forward change

$\Rightarrow$  backward error is the negative of  
the forward error, i.e.,

$$\varepsilon(-h) = -\varepsilon(h)$$

$\Rightarrow \varepsilon(h)$  is an odd function &  
has <sup>only</sup> odd powers in Taylor expansion.

For overall error, we compute  
error on a single step times the  
# of steps.

(9)

So if the time interval of interest is  $\Delta$ , then # of steps is  $\Delta/h \Rightarrow$

$$\text{total error is } \varepsilon(h) \cdot \frac{\Delta}{h}$$

i.e., an even function of error.

w/ first term proportional to  $h^2$ .

Slight catch: For the 1st step of leap frog, we take a  $1/2$  step using Euler method & introduce an error of order  $h^2$  (as is case w/ Euler method).

However, Euler method higher order terms are not even.

So the total error has even & odd powers.

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Can solve this problem using  
the modified midpoint method.

Suppose we want to solve starting

@  $t$  & ending @  $t+h$  where  
 $h$  is not  
small,

use  $n$  steps of size  $h = H/n$

can write leap frog as

$$x_0 = x(t)$$

$$y_1 = x_0 + \frac{1}{2}h f(x_0, t)$$

& then

$$x_1 = x_0 + h f(y_1, t + \frac{1}{2}h)$$

$$y_2 = y_1 + h f(x_1, t + h)$$

$$x_2 = x_1 + h f(y_2, t + \frac{3}{2}h)$$

etc. general form is

$$y_{m+1} = y_m + h f(x_m, t + mh)$$

$$x_{m+1} = x_m + h f(y_{m+1}, t + (m + \frac{1}{2})h)$$

last 2 points are

$$y_n = x(t+h - \frac{1}{2}h)$$

$$x_n = x(t+h)$$

usually would take  $x_n$  as solution but we can also estimate  $x(t+h)$  using Euler method as

$$x(t+h) = y_n + \frac{1}{2}h f(x_n, t+h)$$

can take the average of the

two estimates for  $x(t+h)$  to be the estimate:

$$x(t+h) = \frac{1}{2} [x_n + y_n + \frac{1}{2}h f(x_n, t+h)]$$

Can show that doing so

cancels out the odd order error terms introduced in the

1st step of Euler method

(need to track errors carefully to see this)

contains only even-order error terms. Remember how this was useful for Romberg integration?

This will be useful for a similar method for solving diff. eq's

### Burkirsch - Stoer Method

reminiscent of Romberg integration

Suppose we wish to solve

$$\frac{dx}{dt} = f(x,t)$$

from  $t$  to  $t+H$

Begin by using a single step of size  $H$  & use modified midpoint method

Let  $h_1 = H$

gives an estimate of  $x(t+H)$ ,

call it  $R_{1,1}$ .

Now go back to time  $t$ ,

divide interval into 2 steps of

size  $h_2 = \frac{1}{2}H$ . Gives

another estimate of  $x(t+H)$ ,

call it  $R_{2,1}$ .

Since total error of modified midpoint method is an odd function of step size, we have that

$$x(t+H) = R_{2,1} + c_1 h_2^2 + O(h_2^4)$$

where  $c_1$  is some constant

(14)

Also consider that

$$x(t+h) = R_{1,1} + c_1 h_1^2 + O(h_1^4)$$

$$= R_{1,1} + 4c_1 h_2^2 + O(h_2^4)$$

using  $h_1 = 2h_2$

Then since both of these are equal

$\rightarrow x(t+h)$ , equate them

~~the last term.~~ to  
get

$$c_1 h_2^2 = \frac{1}{3} (R_{2,1} - R_{1,1}) + O(h_2^4)$$

Substitute into the above to get

$$\begin{aligned} x(t+h) &= R_{2,1} + \frac{1}{3} (R_{2,1} - R_{1,1}) \\ &\quad + O(h_2^4) \end{aligned}$$

- method has ~~error of order~~  $h_2^4$

- call the new estimate

$$R_{2,2} = R_{2,1} + \frac{1}{3} (R_{2,1} - R_{1,1})$$

can continue w/ this idea

- increase # of steps to 3

$$\text{so that } h_3 = \frac{h}{3}$$

- solve from  $t$  to  $t+h$  to get  
estimate  $R_{3,1}$

calculate estimate

$$R_{3,2} = R_{3,1} + \frac{4}{5} (R_{3,1} - R_{2,1})$$

(reasoning same as before)

we can then write

$$x(t+h) = R_{3,2} + c_2 h_3^4 + O(h_3^6)$$

w/  $c_2$  a constant

From before we have that

$$x(t+h) = R_{2,2} + c_2 h_2^4 + O(h_2^6)$$

$$= R_{2,2} + \frac{81}{16} c_2 h_3^4 + O(h_2^6)$$

$$\text{using } h_2 = \frac{3}{2} h_3$$

(16)

Equating & rearranging gives

$$C_2 h_3^4 = \frac{16}{65} (R_{3,2} - R_{2,2})$$

substituting back in gives

$$x(t+H) = R_{3,3} + O(h_3^6)$$

$$\text{where } R_{3,3} = R_{3,2} + \frac{16}{65} (R_{3,2} - R_{2,2})$$

result is accurate to order  $h_3^6$

w/ just 3 steps.

Power is the cancellation of higher order errors on successive steps

doing w/ <sup>modified</sup> midpoint method having only even-order terms,

can continue this process iteratively & estimate B

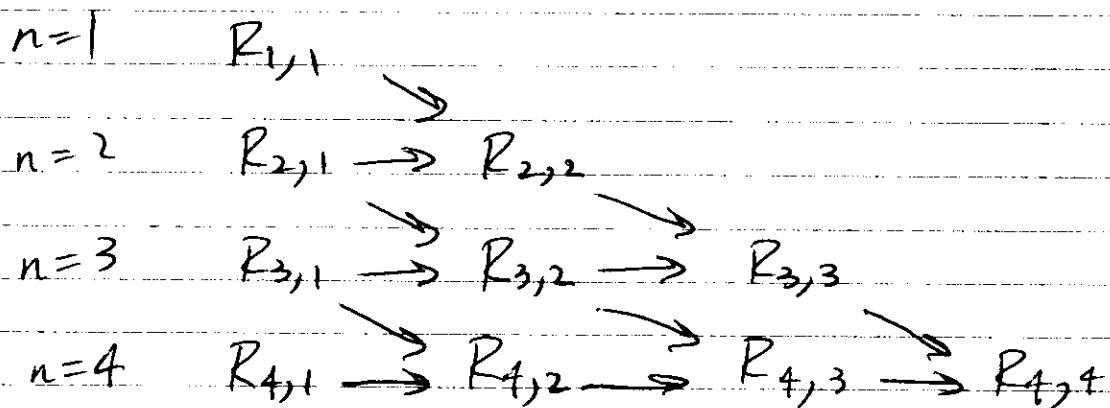
$$x(t+H) = R_{n,m+1} + O(h_n^{2m+2})$$

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w/  $m < n$  +

$$R_{n,m+1} = R_{n,m} + \frac{R_{n,m} - R_{n-1,m}}{[n/(n-1)]^{2m} - 1}$$

diagram similar to that for Romberg integration



modified  
midpoint  
method

Richardson extrapolation

also get estimate of error as

$$\frac{R_{n,m} - R_{n-1,m}}{[n/(n-1)]^{2m} - 1}$$

Limitations :

- 1) answer only accurate for  $x(t+H)$
- 2) method only converges quickly if the power series expansion of  $x(t+H)$  converges quickly.

Can overcome these problems by dividing interval into smaller ones

+ apply ~~Bulirsch~~ - Stoer to Bulirsch

each one.